

Periodic solutions of Schrodinger equation in Hilbert space.

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Necessary and sufficient conditions for existence of boundary value problem of Schrodinger equation are obtained in linear and nonlinear cases. Periodic analytical solutions are represented using generalized Green's operator.

### **Auxiliary result(Linear case).**

**Statement of the problem.** Consider the next boundary value problem for Shrodinger equation

$$\frac{d\varphi(t)}{dt} = -iH_0\varphi(t) + f(t), t \in [0; w] \quad (1)$$

$$\varphi(0) - \varphi(w) = \alpha \in D \quad (2)$$

in a Hilbert space  $H_T$ , where, for each  $t \in [0; w]$ , the unbounded operator  $H_0$  has the form [1]

$$H_0 = i \begin{pmatrix} 0 & T \\ -T & 0 \end{pmatrix} = \begin{pmatrix} T & 0 \\ 0 & T \end{pmatrix} \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix} = i \begin{pmatrix} 0 & I \\ -I & 0 \end{pmatrix} \begin{pmatrix} T & 0 \\ 0 & T \end{pmatrix},$$

for simplicity. In more general case operator  $H_0$  has the next form

$$H_0 = iJ \begin{pmatrix} T & 0 \\ 0 & T \end{pmatrix} = i \begin{pmatrix} T & 0 \\ 0 & T \end{pmatrix} J, \quad J = J^* = J^{-1},$$

where  $T$  is strongly positive self-adjoint operator in the Hilbert space  $H$ . Since operator  $T^2$  is closed, then domain  $D(T)$  of operator  $T$  is Hilbert space with scalar product  $(Tu, Tu)$ . The space  $H_T = H \oplus H$  and operator  $H_0$  is self-adjoint on domain  $D = D(T) \oplus D(T)$  with product

$$(< u, v >, < u, v >)_{H^T} = (Tu, Tu)_H + (Tv, Tv)_H$$

and infinitesimal generator of strongly continuous evolution semigroup

$$U(t) := U(t, 0) = \begin{pmatrix} \cos tT & \sin tT \\ -\sin tT & \cos tT \end{pmatrix}, \quad U^n(t) = \begin{pmatrix} \cos ntT^{\frac{1}{2}} & \sin ntT^{\frac{1}{2}} \\ -\sin ntT^{\frac{1}{2}} & \cos ntT^{\frac{1}{2}} \end{pmatrix} = U(nt),$$

$\|U^n(t)\| = 1, n \in \mathbb{N}$  (nonexpanding group);  $\varphi(t) = (\varphi_1(t), \varphi_2(t))^T$ ,  $\alpha = (\alpha_1, \alpha_2)^T$ ,  $f(t) = (f_1(t), f_2(t))^T$ . Solutions of equation (1) can be represented in the next form

$$\varphi(t) = U(t)c + \int_0^t U(t)U^{-1}(\tau)f(\tau)d\tau,$$

for any element  $c \in H_T$ . Substitute in condition (2) we obtain that solvability of boundary value problem (1), (2) is equivalent solvability the next operator equation

$$(I - U(w))c = g, \quad (3)$$

where  $g = \alpha + U(w) \int_0^w U^{-1}(\tau)f(\tau)d\tau$ . Consider the case when the set of values of  $I - U(w)$  is closed  $R(I - U(w)) = \overline{R(I - U(w))}$ . Since  $\|U^n(w)\| = \|U(w)\| = 1$  for all  $n \in \mathbb{N}$  then [2] the operator system (3) is solvable if and only if

$$U_0(w)g = 0, \quad (4)$$

where

$$U_0(w) = \lim_{n \rightarrow \infty} \frac{\sum_{k=0}^n U^k(w)}{n} = \lim_{n \rightarrow \infty} \frac{\sum_{k=0}^n U(kw)}{n} -$$

orthoprojector, which projects the space  $H_T$  onto subspace  $1 \in \sigma(U(w))$ . Under this condition solutions of (3) have the form

$$c = U_0(w)\bar{c} + \left( \sum_{k=0}^{\infty} (\mu - 1)^k \left\{ \sum_{l=0}^{\infty} \mu^{-l-1} (U(w) - U_0(w))^l \right\}^{k+1} - U_0(w) \right) g,$$

for  $\mu > 1, |1 - \mu| < \frac{1}{\|R_\mu(U(w))\|}$  and any  $\bar{c} \in H_T$ . Then we can formulate first result as lemma.

**Lemma 1.** *Let the operator  $I - U(w)$  has a closed image  $R(I - U(w)) = \overline{R(I - U(w))}$ .*

1. *There exist solutions of boundary value problem (1), (2) if and only if*

$$U_0(w)(\alpha + \int_0^w U^{-1}(\tau)f(\tau)d\tau) = 0. \quad (5)$$

2. *Under condition (5), solutions of (1), (2) have the form*

$$\varphi(t, \bar{c}) = U(t)U_0(w)\bar{c} + (G[f, \alpha])(t), \quad (6)$$

where

$$(G[f, \alpha])(t) = U(t) \sum_{k=0}^{\infty} (\mu - 1)^k \left\{ \sum_{l=0}^{\infty} \mu^{-l-1} (U(w) - U_0(w))^l \right\}^{k+1} (\alpha + \int_0^w U(w)U^{-1}(\tau)f(\tau)d\tau) -$$

$$-U(t)U_0(w)(\alpha + \int_0^w U(w)U^{-1}(\tau)f(\tau)d\tau) + \int_0^t U(t)U^{-1}(\tau)f(\tau)d\tau,$$

is the generalized Green operator of the boundary value problem (1), (2) for  $\mu > 1, |1-\mu| < 1/||R_\mu(U(w))||$ .

Now we show that condition  $R(I - U(w)) = \overline{R(I - U(w))}$  of lemma 1 can be omitted and in the different senses boundary value problem (1), (2) is always resolvable.

1) Classical generalized solutions.

Consider case when the set of values of  $I - U(w)$  is closed ( $R(I - U(w)) = \overline{R(I - U(w))}$ ). Then [3]  $g \in R(I - U(w))$  if and only if  $\mathcal{P}_{N(I - U(w))^*}g = 0$  and the set of solutions of (3) has the form [3]  $c = G[g] + U_0(w)\bar{c}, \forall \bar{c} \in H_T$ , where [2] and [3]

$$G[g] = (I - U(w))^+g = ((I - (U(w) - U_0(w))^{-1} - U_0(w))g$$

is generalized Green operator (or in the form of convergent series).

2) Strong generalized solutions. Consider the case when  $R(I - U(w)) \neq \overline{R(I - U(w))}$  and  $g \in \overline{R(I - U(w))}$ . We show that operator  $I - U(w)$  may be extended to  $\overline{I - U(w)}$  in such way that  $R(\overline{I - U(w)})$  is closed.

Since the operator  $I - U(w)$  is bounded the next representation of  $H_T$  in the direct sum is true

$$H_T = N(I - U(w)) \oplus X, H_T = \overline{R(I - U(w))} \oplus Y,$$

with  $X = N(I - U(w))^\perp = \overline{R(I - U(w))}$  and  $Y = \overline{R(I - U(w))}^\perp = N(I - U(w))$ . Let  $E = H_T/N(I - U(w))$  is quotient space of  $H_T$ ,  $\mathcal{P}_{\overline{R(I - U(w))}}$  and  $\mathcal{P}_{N(I - U(w))}$  are orthoprojectors, which project onto  $\overline{R(I - U(w))}$  and  $N(I - U(w))$  respectively. Then operator

$$\mathcal{I} - \mathcal{U}(w) = \mathcal{P}_{\overline{R(I - U(w))}}(I - U(w))j^{-1}p : X \rightarrow R(I - U(w)) \subset \overline{R(I - U(w))},$$

is linear, continuous and injective. Here

$$p : X \rightarrow E = H_T/N(I - U(w)), \quad j : H_T \rightarrow E$$

are continuous bijection and projection respectively. The triple  $(H_T, E, j)$  is a locally trivial bundle with typical fiber  $H_1 = \mathcal{P}_{N(I - U(w))}H$  [4]. In this case [5, p.26,29] we can define strong generalized solution of equation

$$(\mathcal{I} - \mathcal{U}(w))x = g, x \in X. \tag{7}$$

Fill up the space  $X$  in the norm  $\|x\|_{\overline{X}} = \|(\mathcal{I} - \mathcal{U}(w))x\|_F$ , where  $F = \overline{R(I - U(w))}$  [5]. Then extended operator

$$\overline{\mathcal{I} - \mathcal{U}(w)} : \overline{X} \rightarrow \overline{R(I - U(w))}, X \subset \overline{X}$$

is homeomorphism of  $\overline{X}$  and  $\overline{R(I - U(w))}$ . By virtue of construction of strong generalized solution [5] equation

$$(\overline{\mathcal{I} - \mathcal{U}(w)})\bar{\xi} = g,$$

has a unique solution  $(\overline{\mathcal{I} - \mathcal{U}(w)})^{-1}g$  which is called generalized solution of equation (7).

**Remark 1.** It should be noted that there are exists next extensions of spaces and corresponding operators

$$\overline{p} : \overline{X} \rightarrow \overline{E}, \quad \overline{j} : \overline{H}_T \rightarrow \overline{E}, \quad \overline{\mathcal{P}}_X = \mathcal{P}_{\overline{X}} : \overline{H}_T \rightarrow \overline{X}, \quad \overline{G} : \overline{R(I - U(w))} \rightarrow \overline{X},$$

where

$$\overline{H}_T = N(I - U(w)) \oplus \overline{X}; \quad \overline{p}(x) = p(x), x \in X; \quad \overline{j}(x) = j(x), x \in H_T,$$

$$\overline{\mathcal{P}}_X(x) = \mathcal{P}_X(x), x \in H_T \quad (\mathcal{P}_X = \mathcal{P}_X^2 = \mathcal{P}_X^*); \quad \overline{G}[g] = G[g], g \in R(I - U(w)).$$

Then the operator  $\overline{I - U(w)} = (\overline{\mathcal{I} - \mathcal{U}(w)})\mathcal{P}_{\overline{X}} : \overline{H}_T \rightarrow \overline{H}_T$  is extension of  $I - U(w)$ ,  $(\overline{I - U(w)})c = (I - U(w))c$  for all  $c \in H_T$ .

3) Strong pseudosolutions.

Consider element  $g \notin \overline{R(I - U(w))}$ . This condition is equivalent  $\mathcal{P}_{N(I - U(w))^*}g \neq 0$ . In this case there are exists elements from  $\overline{H}_T$  which minimise norm  $\|(\overline{I - U(w)})\xi - g\|_{\overline{H}_T}$  :

$$\xi = (\overline{\mathcal{I} - \mathcal{U}(w)})^{-1}g + \mathcal{P}_{N(I - U(w))}\bar{c}, \forall \bar{c} \in H_T.$$

These elements we call *strong pseudosolutions* by analogy of [3].

Now we formulate the full theorem of solvability.

**Theorem 1.** *Boundary value problem (1), (2) is always resolvable.*

1) a) *There are exists classical or strong generalized solutions of (1), (2) if and only if*

$$U_0(w)(\alpha + \int_0^w U^{-1}(\tau)f(\tau)d\tau) = 0. \quad (8)$$

*If  $(\alpha + \int_0^w U^{-1}(\tau)f(\tau)d\tau) \in R(I - U(w))$  then solutions of (1), (2) will be classical.*

b) *Under assumption (8) solutions of (1), (2) have the form*

$$\varphi(t, \bar{c}) = U(t)U_0(w)\bar{c} + (\overline{G[f, \alpha]})(t),$$

where  $(\overline{G[f, \alpha]})(t)$  - is extension of operator  $(G[f, \alpha])(t)$ ;

3) a) There are exists strong pseudosolutions if and only if

$$U_0(w)(\alpha + \int_0^w U^{-1}(\tau)f(\tau)d\tau) \neq 0. \quad (9)$$

b) Under assumption (9) strong pseudosolutions of (1), (2) have the form

$$\varphi(t, \bar{c}) = U(t)U_0(w)\bar{c} + (\overline{G[f, \alpha]})(t),$$

where

$$(\overline{G[f, \alpha]})(t) = U(t)\overline{G}[g] + \int_0^t U(t)U^{-1}(\tau)f(\tau)d\tau = U(t)(\overline{\mathcal{I} - \mathcal{U}(w)})^{-1}g + \int_0^t U(t)U^{-1}(\tau)f(\tau)d\tau.$$

### **Main result (Nonlinear case). Generalization of Lyapunov-Schmidt method**

In the Hilbert space  $H_T$  defined below we consider the boundary value problem

$$\frac{d\varphi(t)}{dt} = -iH_0\varphi(t) + \varepsilon Z(\varphi(t), t, \varepsilon) + f(t), \quad (10)$$

$$\varphi(0, \varepsilon) - \varphi(w, \varepsilon) = \alpha. \quad (11)$$

We seek a bounded solution  $\varphi(t, \varepsilon)$  of boundary value problem (10), (11) that becomes one of the solutions of the generating equation (1), (2)  $\varphi_0(t, \bar{c})$  in the form (6) for  $\varepsilon = 0$ .

To find a necessary condition on the operator function  $Z(\varphi, t, \varepsilon)$ , we impose the joint constraints

$$Z(\cdot, \cdot, \cdot) \in C([0; w], H_T) \times C[0, \varepsilon_0] \times C[\|\varphi - \varphi_0\| \leq q],$$

where  $q$  is some positive constant.

The main idea of the next results is presented in [6] for investigating of bounded solutions.

Let us show that this problem can be solved with the use of the operator equation for generating amplitudes

$$F(\bar{c}) = U_0(w) \int_0^w U^{-1}(\tau)Z(\varphi_0(\tau, \bar{c}), \tau, 0)d\tau = 0. \quad (12)$$

**Theorem 2.** (necessary condition) *Let the nonlinear boundary value problem (10), (11) has a bounded solution  $\varphi(\cdot, \varepsilon)$  that becomes one of the solutions  $\varphi_0(t, \bar{c})$  of the generating equation (1), (2) with constant  $\bar{c} = c^0$ ,  $\varphi(t, 0) = \varphi_0(t, c^0)$  for  $\varepsilon = 0$ . Then this constant should satisfy the equation for generating amplitudes (12).*

To find a sufficient condition for the existence of solutions of boundary value problem (10), (11) we additionally assume that the operator function  $Z(\varphi, t, \varepsilon)$  is strongly differentiable in a neighborhood of the generating solution ( $Z(\cdot, t, \varepsilon) \in C^1[||\varphi - \varphi_0|| \leq q]$ ).

This problem can be solved with the use of the operator

$$B_0 = \frac{dF(\bar{c})}{d\bar{c}}|_{\bar{c}=c^0} = \int_{-\infty}^{+\infty} H(t)A_1(t)T(t, 0)P_+(0)\mathcal{P}_{N(D)}dt : H \rightarrow H,$$

where  $A_1(t) = Z^1(v, t, \varepsilon)|_{v=\varphi_0, \varepsilon=0}$  (the Fréchet derivative).

**Theorem 3.** (sufficient condition) *Let the operator  $B_0$  satisfy the following conditions:*

- 1) *The operator  $B_0$  is Moore-Penrouse pseudoinvertible;*
- 2)  *$\mathcal{P}_{N(B_0^*)}U(w) = 0$ .*

*Then for arbitrary element  $c = c^0 \in H_T$ , satisfying the equation for generating amplitudes (12), there exists at least one solution of (10), (11).*

*This solution can be found with the use of the iterative process:*

$$\begin{aligned}\bar{v}_{k+1}(t, \varepsilon) &= \varepsilon G[Z(\varphi_0(\tau, c^0) + v_k, \tau, \varepsilon), \alpha](t), \\ c_k &= -B_0^+U_0(w) \int_0^w U^{-1}(\tau)\{A_1(\tau)\bar{v}_k(\tau, \varepsilon) + \mathcal{R}(v_k(\tau, \varepsilon), \tau, \varepsilon)\}d\tau, \\ v_{k+1}(t, \varepsilon) &= U(t)U_0(w)c_k + \bar{v}_{k+1}(t, \varepsilon),\end{aligned}$$

$$\varphi_k(t, \varepsilon) = \varphi_0(t, c^0) + v_k(t, \varepsilon), k = 0, 1, 2, \dots, \quad v_0(t, \varepsilon) = 0, \varphi(t, \varepsilon) = \lim_{k \rightarrow \infty} \varphi_k(t, \varepsilon).$$

**Remark 2.** Proof of theorems 2 and 3 follows directly from works [6], [7].

### Relationship between necessary and sufficient conditions.

First, we formulate the following assertion.

**Corollary.** *Let a functional  $F(\bar{c})$  have the Fréchet derivative  $F^{(1)}(\bar{c})$  for each element  $c^0$  of the Hilbert space  $H$  satisfying the equation for generating constants (12). If  $F^{(1)}(\bar{c})$  has a bounded inverse, then boundary value problem (10), (11) has a unique solution for each  $c^0$ .*

**Remark 3.** If assumptions of the corollary are satisfied, then it follows from its proof that the operators  $B_0$  and  $F^{(1)}(c^0)$  are equal. Since the operator  $F^{(1)}(\bar{c})$  is invertible, it

follows that assumptions 1 and 2 of Theorem 3 are necessarily satisfied for the operator  $B_0$ . In this case, boundary value problem (10), (11) has a unique bounded solution for each  $c^0 \in H_T$  satisfying (12). Therefore, the invertibility condition for the operator  $F^1(\bar{c})$  relates the necessary and sufficient conditions. In the finite-dimensional case, the condition of invertibility of the operator  $F^{(1)}(\bar{c})$  is equivalent to the condition of simplicity of the root  $c^0$  of the equation for generating amplitudes [3].

In such way we generalize the well-known method of Lyapunov-Schmidt. It should be emphasized that theorem 2 and 3 give us condition of chaotic behavior of (10), (11) [8].

**Example.** Now we illustrate obtained assertion. Consider the next differential equation in separable Hilbert space  $H$

$$\ddot{y}(t) + Ty(t) = \varepsilon(1 - \|y(t)\|^2)\dot{y}(t), \quad (13)$$

$$y(0) = y(w), \quad \dot{y}(0) = \dot{y}(w), \quad (14)$$

where  $T$  is unbounded operator with compact  $T^{-1}$ . Then there is exists orthonormal basis  $e_i \in H$  such that  $y(t) = \sum_{i=1}^{\infty} c_i(t)e_i$  and  $Ty(t) = \sum_{i=1}^{\infty} \lambda_i c_i(t)e_i$ ,  $\lambda_i \rightarrow \infty$ . Operator system (10), (11) for boundary value problem (13), (14) in this case will be equivalent the next countable system of ordinary differential equations ( $c_k(t) = x_k(t)$ )

$$\dot{x}_k(t) = \sqrt{\lambda_k}y_k(t), \quad k = 1, 2, \dots,$$

$$\dot{y}_k(t) = -\sqrt{\lambda_k}x_k(t) + \varepsilon\sqrt{\lambda_k}\left(1 - \sum_{j=1}^{\infty} x_j^2(t)\right)y_k(t), \quad (15)$$

$$x_k(0) = x_k(w), y_k(0) = y_k(w). \quad (16)$$

We find solutions of these equations such that for  $\varepsilon = 0$  turns in one of the solutions of generating equation. Consider critical case  $\lambda_i = 4\pi^2 i^2/w^2, i \in \mathbb{N}$ . Let  $w = 2\pi$ . In that case the set of all periodic solutions of (15), (16) have the form

$$x_k(t) = \cos(kt)c_1^k + \sin(kt)c_2^k,$$

$$y_k(t) = -\sin(kt)c_1^k + \cos(kt)c_2^k,$$

for all pairs of constant  $c_1^k, c_2^k \in \mathbb{R}, k \in \mathbb{N}$ . Equation for generating amplitudes (12) in this case will be equivalent the next countable systems of algebraic nonlinear equations

$$(c_1^k)^3 + 2 \sum_{j=1, j \neq k} (c_1^k(c_1^j)^2 + c_1^k(c_2^j)^2) + c_1^k(c_2^k)^2 - 4c_1^k = 0,$$

$$(c_2^k)^3 + 2 \sum_{j=1, j \neq k} (c_2^k(c_1^j)^2 + c_2^k(c_2^j)^2) + (c_1^k)^2 c_2^k - 4c_2^k = 0, k \in \mathbb{N}.$$

Then we can obtain the next result

**Theorem 4**(necessary condition of van der Pol's equation). *Let the boundary value (15), (16) have a bounded solution  $\varphi(\cdot, \varepsilon)$  that becomes one of the solutions of the generating equations with pairs of constant  $(c_1^k, c_2^k), k \in \mathbb{N}$ . Then only finite number of these pairs are not equal zero. Moreover, if  $(c_1^{k_i}, c_2^{k_i}) \neq (0, 0), i = \overline{1, N}$  then these constants lie on  $N$ -dimensional torus in infinite dimensional space of constants*

$$(c_1^{k_i})^2 + (c_2^{k_i})^2 = \left(\frac{2}{\sqrt{2N-1}}\right)^2, i = \overline{1, N}.$$



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